

## 2. Fundament. závislosti: $I = I(\phi)$ & $V = V(\phi)$

1934 greenshields (1935 elabek)

$$\alpha_k(\tau) = v_k(\tau) = \lambda \frac{v_{k-1}(\tilde{\tau}) - v_k(\tau)}{[x_{k-1}(\tilde{\tau}) - x_k(\tau)]^2} \quad \begin{matrix} \text{maly headway!} \\ \leftarrow N, C \end{matrix}$$

$\lambda > 0$

reakční doba

$$\tilde{\tau} = \tau + \tilde{\Delta\tau}$$

$$v'_k(\tau) = -\lambda \frac{d}{d\tau} \frac{1}{x_{k-1}(\tilde{\tau}) - x_k(\tau)}$$

$$v'_{k-1}(\tilde{\tau}) = v_k(\tau)$$

$$v_k(\tau) = \frac{-\lambda}{x_{k-1}(\tilde{\tau}) - x_k(\tilde{\tau})} + C_k$$

$$\bar{\rho}_k^{-1}(\tilde{\tau}) := x_{k-1}(\tilde{\tau}) - x_k(\tilde{\tau}) \quad k = 1, 2, \dots, N \quad \} \Rightarrow \bar{\rho}_k = \frac{1}{n_k}$$

$\bar{\rho}_k(\tilde{\tau}) \quad \text{tjv. individualní hustoty}$

$\left. \begin{matrix} \bar{\rho} = \frac{1}{N} \sum_{k=1}^N \bar{\rho}_k \\ \text{teoretická hustota: } \bar{\rho} = \mathbb{E}\left(\frac{1}{n}\right) \end{matrix} \right\} \quad \begin{matrix} \text{m}_k(\tilde{\tau}) \dots \text{průměrný rozestup} \\ (\bar{\rho} \text{ nemá konkrétní hustota}) \\ \bar{\rho} = \text{aritm. průměr individuálních hustot} \end{matrix}$

$$v_k(\tau) = \frac{-\lambda}{x_{k-1}(\tilde{\tau}) - x_k(\tilde{\tau})} + C_k \quad \left/ \frac{1}{N} \sum_k \right.$$

$$v_k(\tau) = -\lambda \bar{\rho}_k + C_k$$

$$V = -\lambda \bar{\rho} + C$$

Běžné stabilitu zajišťují podmínky :  $\bar{\rho} = 0 \Rightarrow V = V_0$  → optimálnost

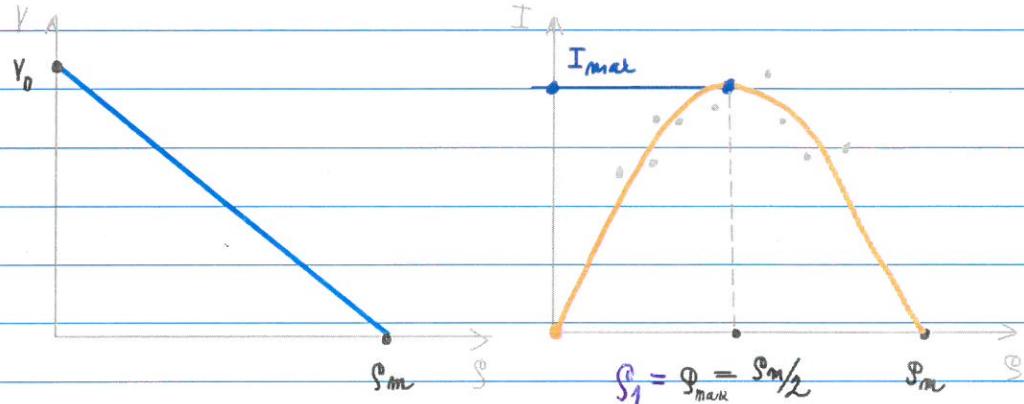
$\bar{\rho} = \bar{\rho}_m \Rightarrow V = 0$

$C = V_0 \quad \& \quad \lambda = \frac{V_0}{S_m}$

$$I = g V$$

$$V = V_0 \left(1 - \frac{g}{g_m}\right) \Rightarrow I = V_0 g \left(1 - \frac{g}{g_m}\right)$$

greenshield theory



$$I_{max} = \frac{N}{\tau_N - \tau_1} = \frac{1}{\langle t \rangle} \quad \langle t \rangle \equiv E(t) = \int_R^{\infty} t f(t) dt = \frac{3}{2} \text{ [v sekundach]}$$

$$\int_R^{\infty} x^2 e^{-ax} dx = \frac{2}{a}$$

$$\Rightarrow I_{max} = \frac{2}{3} \text{ rr/s} = \frac{2}{3} VR \cdot 3600 h^{-1} = 2400 \text{ rr/h}$$

+∞

$$E(\frac{1}{r}) = \int_R^{\infty} \frac{1}{r} \cdot g(r) dr = \int_0^{\infty} 40^2 \cdot e^{-40r} dr = 40 \text{ VR/km}$$

$$g_1 = \frac{g_m}{2} = 40 \text{ rr/km} \Rightarrow g_m = 80 \text{ rr/km}$$

$$g_m = 40 \text{ & } I_{max} = 2400 \text{ & } I = g V \Rightarrow V_{max} = 60 \text{ km/h}$$

$$\Rightarrow V_0 = 120 \text{ km/h}$$

$$V = 120 \left(1 - \frac{g}{80}\right) \quad \& \quad I = 120 g \left(1 - \frac{g}{80}\right)$$

$$\text{Klasika: } \int_{\mathbb{R}} e^{-ax^2} dx = \sqrt{\frac{\pi}{a}} / \frac{d}{da}$$

$$-\int_{\mathbb{R}} x^2 e^{-ax^2} dx = -\frac{1}{2} \sqrt{\pi} \frac{1}{a^{3/2}} \quad \& \quad a = \frac{1}{2s^2}$$

$$\int_{\mathbb{R}} (y-x)^2 e^{-\frac{(y-x)^2}{2s^2}} ds = \int_{\mathbb{R}} z^2 e^{-\frac{z^2}{2s^2}} ds = \sqrt{2\pi} s^3 \stackrel{!}{=} A^1$$

$$\Rightarrow A = \frac{1}{\sqrt{2\pi} s^3}$$

$$\varrho(x, t) = \frac{\partial N}{\partial x} = \frac{1}{\sqrt{2\pi} s^3} \sum_{k=1}^N (x - \alpha_k)^2 e^{-\frac{(x - \alpha_k)^2}{2s^2}}$$

$$\lim_{s \rightarrow 0_+} \varrho(x, t) \stackrel{?}{=} ?$$

$$\left( \lim_{s \rightarrow 0_+} \varrho(x, t), \varphi(x) \right) \stackrel{\text{def.}}{=} \lim_{s \rightarrow 0_+} (\varrho(x, t), \varphi(x)) = \left\| \begin{array}{l} \varrho(x, t) \text{ je regulärne} \\ \text{distributio a } \varphi(x) \in \mathcal{D} \end{array} \right\| =$$

$$= \lim_{s \rightarrow 0_+} \int_{\mathbb{R}} \frac{1}{\sqrt{2\pi} s^3} \sum_{k=1}^N (x - \alpha_k)^2 e^{-\frac{(x - \alpha_k)^2}{2s^2}} \varphi(x) dx =$$

$$= \sum_{k=1}^N \lim_{s \rightarrow 0_+} \frac{1}{\sqrt{2\pi} s^3} \int_{\mathbb{R}} (x - \alpha_k)^2 e^{-\frac{(x - \alpha_k)^2}{2s^2}} \varphi(x) dx = \left\| \begin{array}{l} \text{jakhe ale zamenit} \\ \text{mee!} \end{array} \right\| =$$

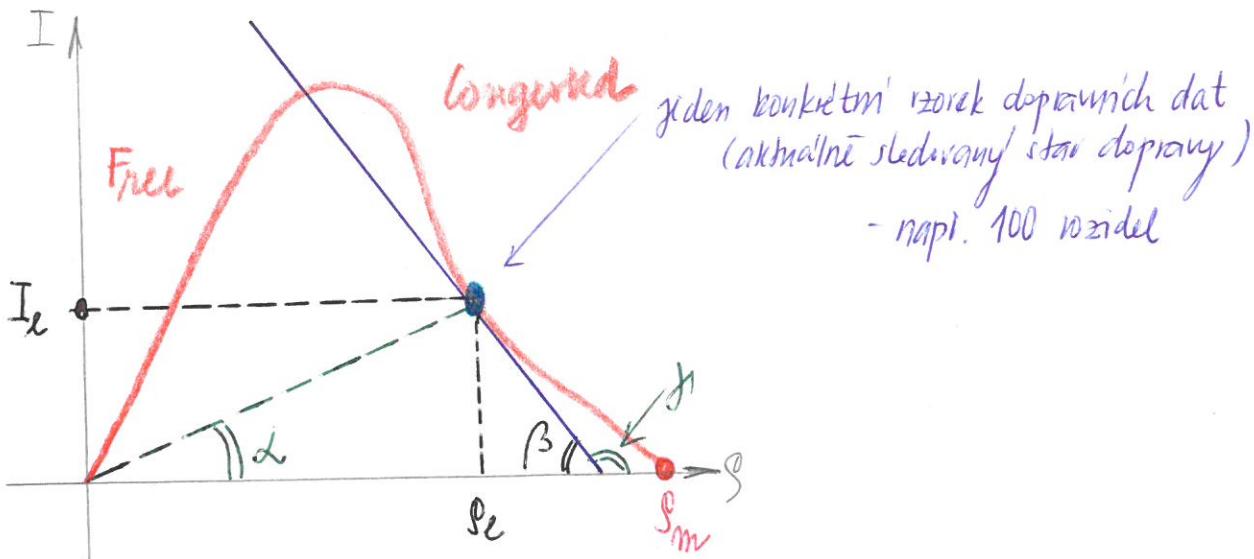
$$= \left\| \begin{array}{l} y = \frac{x - \alpha_k}{s} \\ dy = \frac{dx}{s} \end{array} \right\| = \sum_{k=1}^N \lim_{s \rightarrow 0_+} \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} y^2 e^{-\frac{y^2}{2}} \varphi(y \cdot s + \alpha_k) dy =$$

$$= \left\| \begin{array}{l} |y^2 e^{-\frac{y^2}{2}} \varphi(y \cdot s + \alpha_k)| \leq K \cdot y^2 e^{-\frac{y^2}{2}} \in L(\mathbb{R}) \\ \text{z vlastnost' to dy} \end{array} \right\| =$$

$$= \sum_{k=1}^N \varphi(\alpha_k) \cdot \int_{\mathbb{R}} y^2 e^{-\frac{y^2}{2}} dy \cdot \frac{1}{\sqrt{2\pi}} = \sum_{k=1}^N \varphi(\alpha_k) \cdot \sqrt{2\pi} \cdot \frac{1}{\sqrt{2\pi}} = \sum_{k=1}^N \varphi(\alpha_k) =$$

$$= \left( \sum_{k=1}^N \delta_{\alpha_k}(x), \varphi(x) \right) = \left( \sum_{k=1}^N \delta_{\alpha_k(x)}(x), \varphi(x) \right) \Rightarrow \begin{aligned} \lim_{s \rightarrow 0_+} \varrho(x, y) &= \\ &= \sum_{k=1}^N \delta(x - \alpha_k(x)) \end{aligned}$$

## Illustrace geometrických vztahů mezi FD a dopravními relacemi



①

$$\operatorname{tg} \alpha = \frac{I_e}{s_e} = v_e \quad , \text{ tj. průměrná rychlosť sledovaného řazku}$$

- počet by individualných rychlosťí měl normální rozdělení  $N(\mu, \sigma^2)$ , pak by  $v_e = E(V) = \mu$

②

Rychlosť kinematické vlny

$$- u(p) = V(p) + p \cdot \frac{dV}{ds}$$

$$\frac{dI}{ds} = \parallel \text{tj. týčna ke grafu na kole} \\ \text{v bodě } (s_e, I_e), \text{ resp. její tangenta} \parallel =$$

$$- \frac{d}{dp} \left( s \cdot V(p) \right) = V + p \frac{dV}{dp} \equiv u(p)$$

$$\Rightarrow u(p) = -1 \cdot \operatorname{tg}(\beta) = \operatorname{tg}(\alpha)$$

$$\Rightarrow \underline{\operatorname{tg}(\alpha) = u(p)}$$

## Kinematické dopravní vlny

- kinematickou vlnou rozumíme stokovou (rážovou) vlnu v hustotním profilu
- reprezentuje pohyb maxima hustoty v prostorové cestě
- rychlosť kinematickej vlny nazývame rýchlosť, ktorú sa poruča (t.j. peak hustotného profilu) sŕí prostredom

## Rámcevej odnosom

- z rovnice kontinuity, hydrodynamické alternatív a rosem PDE

$$\frac{\partial \rho}{\partial t} + \frac{\partial I}{\partial x} = 0$$

$$I = \rho \cdot V$$

$$\text{rovnica } \frac{\partial f}{\partial x} + m \frac{\partial f}{\partial y} = 0$$

a fundamentalných hypotéz  $\rightarrow V = V(\rho)$   
 $I = I(\rho)$

$$\stackrel{\Delta}{=} V'$$

$$\frac{\partial \rho}{\partial t} + \frac{\partial I}{\partial x} = \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} (\rho \cdot V(\rho)) = \frac{\partial \rho}{\partial t} + V(\rho) \frac{\partial \rho}{\partial x} + \rho \cdot \frac{dV}{d\rho} \cdot \frac{\partial \rho}{\partial x} = 0$$

$$\frac{\partial \rho}{\partial t} + \left[ V(\rho) + \rho \frac{dV}{d\rho} \right] \cdot \frac{\partial \rho}{\partial x} = 0$$

$$\text{rovnica } f(x, y) = K(y - \mu x)$$

je rýchlosť  
funkcia f

$\mu(V)$  ... propagáciu rýchlosť kinematickej vlny

## Z jednoduchej vzťahu pre rýchlosť $\mu(V)$

Podmínky: 1)  $V(\rho)$  je lineárna funkcia (podľa Greenshielda 1934-35)

$$\Rightarrow \rho = 0 \Rightarrow V = V_0$$

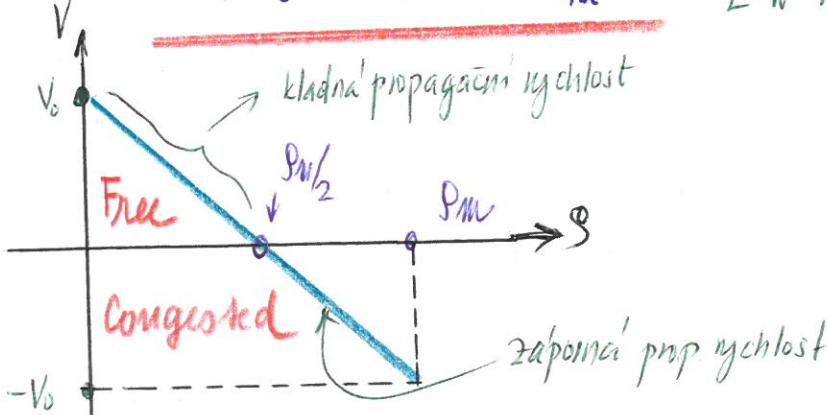
$$\Rightarrow \rho = \rho_m \Rightarrow V = 0$$

$$V = V_0 \left( 1 - \frac{\rho}{\rho_m} \right)$$

$$\Rightarrow V' = \frac{dV}{d\rho} = -\frac{V_0}{\rho_m} \Rightarrow \mu(V) = V_0 - V_0 \frac{\rho}{\rho_m} - \rho \frac{V_0}{\rho_m}$$

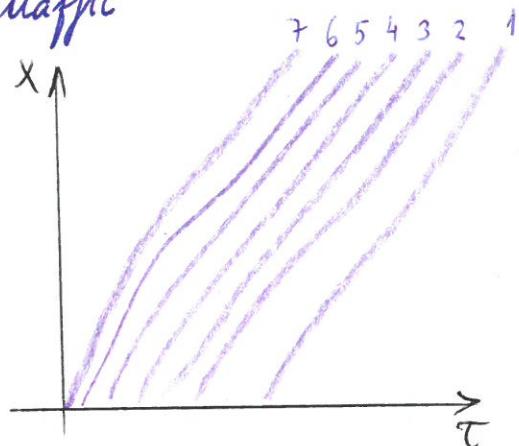
$$\Gamma \quad \mu(V) = V_0 \left( 1 - 2 \frac{\rho}{\rho_m} \right)$$

Bude sa hodit do  
L-W modelu



# Rozbor chování kinematických vln v kontextu třífázové dopravní teorie

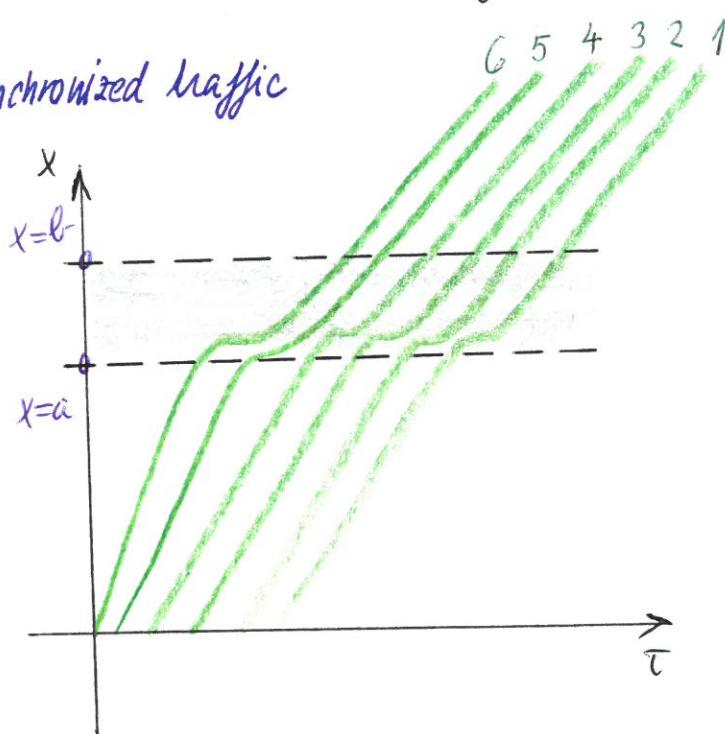
## I. Free traffic



$$\rho \text{ mžke} \& \rho u(\rho) = V_0(1 - 2\frac{\rho}{\rho_m}) \\ \Rightarrow \rho u \doteq V_0$$

- propagací rychlost vlnění splyra's rychlosti vozidel (přiměrnou)

## II. Synchronized traffic

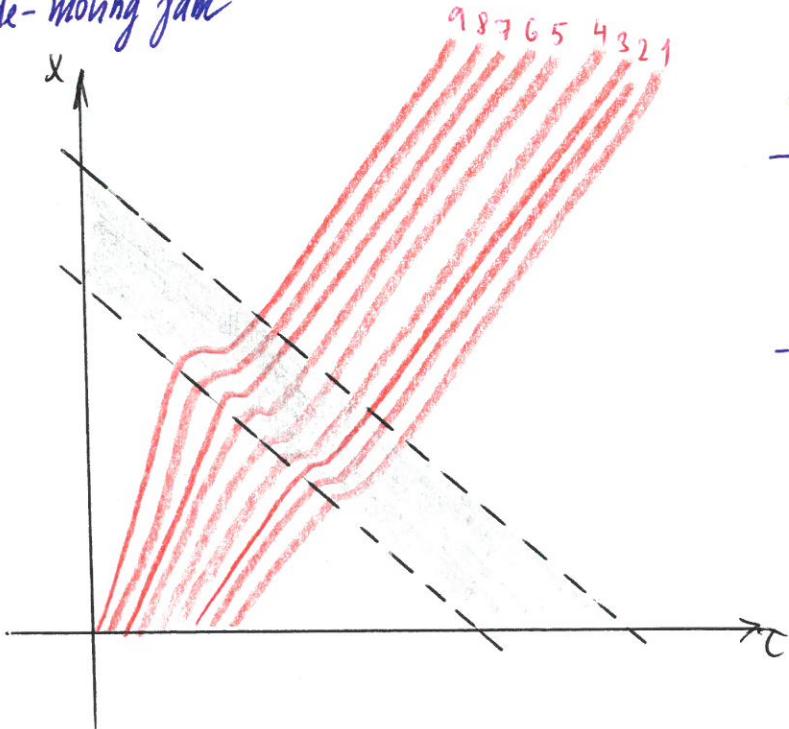


$$\rho \doteq \frac{\rho_m}{2} \Rightarrow \rho u(\rho) \doteq 0$$

nic se nepropaguje

- hustotm' peak zůstává na místě, zde v intervalu  $(a, b)$

## III. Wide-moving jam



$$\rho \gg \frac{\rho_m}{2} \Rightarrow \rho u(\rho) < 0$$

- hustotm' peaky a sunou v opačném směru než je rychlosť vozidel

- uráduje: běžna' rychlosť kongesu:  $\doteq -15 \text{ km/h}$

$$P(\vec{r}) = \frac{1}{Z_N(L)} e^{-\beta \sum_{k=1}^N \varphi(r_k)} \Theta(\vec{r}) \cdot \delta(L - \sum_{k=1}^N r_k) \quad \Theta(r_2, r_3, \dots, r_N)$$

$$\begin{aligned} g(r_1) &= \int_{\mathbb{R}^{N-1}} P(\vec{r}) d(r_2, r_3, \dots, r_N) = \frac{1}{Z_N(L)} \Theta(r_1) e^{-\beta \varphi(r_1)} \int_{\mathbb{R}^{N-1}} e^{-\beta \sum_{k=2}^N \varphi(r_k)} \delta(L - r_1 - \sum_{k=2}^N r_k) d(r_2, \dots, r_N) = \\ &= \Theta(r_1) \frac{\tilde{z}_{N-1}(L-r_1)}{Z_N(L)} e^{-\beta \varphi(r_1)} \Rightarrow g(r) = \Theta(r) \frac{\tilde{z}_{N-1}(L-r)}{Z_N(L)} e^{-\beta \varphi(r)} \end{aligned}$$

$$\varphi(r) = -\ln r \Rightarrow g(r) = \Theta(r) \frac{\tilde{z}_{N-1}(L-r)}{Z_N(L)} r^\beta$$

$$\begin{aligned} \mathcal{L}[Z_N(L)] &= \int_0^\infty \int_{\mathbb{R}^N} \delta(L - \sum_{k=1}^N r_k) \cdot \prod_{k=1}^N \Theta(r_k) \cdot r_k^\beta e^{-\beta L} dr dk dL = |\mathcal{L}[\delta(x-y)] = e^{\delta y x}| = \\ &= \int_{\mathbb{R}^N} \prod_{k=1}^N \Theta(r_k) r_k^\beta e^{-\beta \sum_{k=1}^N r_k} dr = \left( \int_0^\infty \Theta(r) r^\beta e^{-\beta r} dr \right)^N = \mathcal{L}[\Theta(r) \cdot r^\beta] = \\ &= \left( \frac{\Gamma(\beta+1)}{\beta^{\beta+1}} \right)^N = \frac{\Gamma^N(\beta+1)}{\beta^{N\beta+N}} \Rightarrow Z_N(L) = \mathcal{L}^{-1}\left[ \frac{\Gamma^N(\beta+1)}{\beta^{N\beta+N}} \right] \end{aligned}$$

- aproximace v sedlouém bode

$$\frac{d}{ds} \left( \frac{\Gamma^N(\beta+1)}{\beta^{N\beta+N}} e^{\lambda s} \right) = \frac{N}{\beta(\beta+1)} \frac{(N\beta+N-1)!}{s^{N\beta+N-1}} e^{\lambda s} - \frac{N\beta+N}{s^{N\beta+N}} \frac{\lambda e^{\lambda s}}{s^{N\beta+N-1}} = 0$$

$$\lambda = \frac{N(\beta+1)}{L} \stackrel{\Delta}{=} \mu \quad \leftarrow \text{sedlo už má}\right.$$

$$\Rightarrow Z_N(L) = \frac{\Gamma^N(\beta+1)}{\lambda^{N\beta+N}} e^{\lambda L} \quad \& \quad Z_{N-1}(L-r) = \frac{\Gamma^{N-1}(\beta+1)}{\mu^{(N-1)(\beta+1)}} e^{\mu(L-r)}$$

$$\mu = \frac{(N-1)(\beta+1)}{L}$$

Pro  $N=L$  a  $N=N-1$  (tj. pro velký počet částic):

$$\lambda = \mu = \beta + 1$$

$$\Rightarrow Z_N(L) = \left( \frac{\Gamma(\beta+1)}{(\beta+1)^{\beta+1}} \right)^N e^{\lambda N}$$

$$\Rightarrow g(r) = \Theta(r) \frac{(\beta+1)^{\beta+1}}{\Gamma(\beta+1)} r^\beta e^{-(\beta+1)r}$$

(zkrátit k řádu)

## Distribuce světlostí modelu TASEP

- můžeme dátky k za obsahem i-tou funkce



$$P(A|B) = \frac{P(A \cap B)}{P(B)} = \frac{\langle w | (D+E)^{i-1} D E D (D+E)^{N-i} | v \rangle}{\langle w | (D+E)^{i-1} D (D+E)^{N-i} | v \rangle} =$$

$$- \langle w | (D+E)^{i-1} D E^{N-i} | v \rangle$$

$$= p_i(k) \quad D = \frac{1}{\beta} \quad E = \frac{1}{\alpha}$$

$$\alpha + \beta = 1 \Rightarrow p_i(k) = \frac{\alpha \beta^{k-1}}{1 - \beta^{N-i}}$$

$$N \neq \infty \Rightarrow p_i(k) = \alpha \beta^k$$

$$\text{střední hodnota: } \langle k \rangle = \sum_{l=1}^{\infty} l \cdot p_i(l) = \alpha \sum_{l=1}^{\infty} l \cdot \beta^{l-1}$$

$$\sum_{l=1}^{\infty} l \cdot x^{l-1} = s(x) \quad \int s(x) dx = \sum_{l=1}^{\infty} x^l = \frac{x}{1-x}$$

$$\Rightarrow s(x) = \left( \frac{x}{1-x} \right)^1 = \frac{1-x+x}{(1-x)^2} = \frac{1}{(1-x)^2}$$

$$\Rightarrow \langle k \rangle = \alpha \cdot \frac{1}{(1-\beta)^2} = || \alpha + \beta = 1 || =$$

$$= \frac{1}{\alpha}$$

# 5

## Lighthillův-Whithamův dopravní model (1954)

snaha vysvětlit tvar fundamentálního diagramu ✓ ↪ I až zcela neaktuálně

snaha vysvětlit tzv. kinematické vlny (tj. časový vývoj hustotních pulzů)

- řešení' 1974

- 3 zákl. veličiny:  $I, \rho, V$  ( $I = \rho \cdot V$ )

$$I(x, \tau) = I_e(x, \tau) - D \cdot \frac{\partial \rho(x, \tau)}{\partial x} \quad (D > 0)$$

$\uparrow$  efektivní  $\downarrow$  difuzní porucha

$$\Rightarrow V(x, \tau) = \frac{I(x, \tau)}{\rho(x, \tau)} = \underbrace{V_e(x, \tau)}_{\frac{\partial \rho}{\partial x}} - \underbrace{\frac{D}{\rho(x, \tau)} \frac{\partial \rho(x, \tau)}{\partial x}}$$

$\frac{V_e(\rho)}{\rho}$  ← východiska makromodeli

Závazky kontinuity:  $\frac{\partial \rho}{\partial \tau} + \frac{\partial I}{\partial x} = \frac{\partial \rho}{\partial \tau} + \frac{\partial}{\partial x} \left[ \rho \cdot V_e(\rho) - D \frac{\partial \rho}{\partial x} \right] \stackrel{!}{=} 0$

$$\frac{\partial \rho}{\partial \tau} + \frac{\partial \rho}{\partial x} \cdot V_e(\rho) + \rho \frac{dV_e}{d\rho} + D \frac{\partial^2 \rho}{\partial x^2} = 0$$

1.  $\frac{\partial \rho}{\partial \tau} + \frac{\partial \rho}{\partial x} \left[ V_e(\rho) + \rho \frac{dV_e}{d\rho} \right] = D \frac{\partial^2 \rho}{\partial x^2}$

Greenshields:

2.  $V_e(\rho) = V_0 \left( 1 - \frac{\rho}{\rho_m} \right)$

$$\frac{\partial \rho}{\partial \tau} + \frac{\partial \rho}{\partial x} \left[ V_0 - V_0 \frac{\rho}{\rho_m} - \rho \frac{V_0}{\rho_m} \right] = D \frac{\partial^2 \rho}{\partial x^2}$$

$\downarrow$   $V_0 \left( 1 - 2 \frac{\rho(x, \tau)}{\rho_m} \right) = \mu(x, \tau)$  ↪ rychlosť kinematickej vlny

$\mu \bullet$   $\mu(x, \tau) = \frac{\rho_m}{2 V_0} (V_0 - \mu(x, \tau))$

$$\frac{\partial \rho}{\partial \tau} = -\frac{\rho_m}{2 V_0} \frac{\partial \mu}{\partial \tau} \quad \& \quad \frac{\partial \rho}{\partial x} = -\frac{\rho_m}{2 V_0} \frac{\partial \mu}{\partial x} \quad \& \quad \frac{\partial^2 \rho}{\partial x^2} = -\frac{\rho_m}{2 V_0} \frac{\partial^2 \mu}{\partial x^2}$$

$$\frac{\partial \mu}{\partial \tau} + \underline{\mu(x, \tau)} \frac{\partial \mu}{\partial x} = D \frac{\partial^2 \mu}{\partial x^2}$$

konzervovaný tvar PDE

↓  
Burgersova PDE (nelineární)

$$\mu(x, \tau) = - \frac{2D}{\underline{\psi(x, \tau)}} \frac{\partial \psi}{\partial x} \quad \text{Cte - konst}$$

$$1. / \frac{\partial \mu}{\partial \tau} = \frac{2D}{\psi^2(x, \tau)} \frac{\partial \psi}{\partial \tau} \frac{\partial \psi}{\partial x} - \frac{2D}{\psi(x, \tau)} \frac{\partial^2 \psi}{\partial x \partial \tau}$$

$$\mu \cdot / \frac{\partial \mu}{\partial x} = \frac{2D}{\psi^2(x, \tau)} \left( \frac{\partial \psi}{\partial x} \right)^2 - \frac{2D}{\psi(x, \tau)} \frac{\partial^2 \psi}{\partial x^2}$$

$$D \cdot / \frac{\partial^2 \mu}{\partial x^2} = - \frac{4D}{\psi^3(x, \tau)} \left( \frac{\partial \psi}{\partial x} \right)^3 + \underbrace{\frac{4D}{\psi^2(x, \tau)} \frac{\partial \psi}{\partial x} \cdot \frac{\partial^2 \psi}{\partial x^2}}_{6D} + \underbrace{\frac{2D}{\psi^2(x, \tau)} \frac{\partial \psi}{\partial x} \frac{\partial^2 \psi}{\partial x^2}}_{-2D} - \frac{2D}{\psi(x, \tau)} \frac{\partial^3 \psi}{\partial x^3}$$

$$\psi(x, \tau) \frac{\partial \psi}{\partial \tau} \cdot \frac{\partial \psi}{\partial x} - \psi^2(x, \tau) \frac{\partial^2 \psi}{\partial x \partial \tau} - 2D \frac{\partial \psi}{\partial x} \left( \frac{\partial \psi}{\partial x} \right)^2 + 2D \frac{\partial \psi}{\partial x} \mu(x, \tau) \cdot \frac{\partial^2 \psi}{\partial x^2} = \\ = -2D \cancel{\left( \frac{\partial \psi}{\partial x} \right)^3} + 3D \psi(x, \tau) \frac{\partial \psi}{\partial x} \frac{\partial^2 \psi}{\partial x^2} - D \psi^2(x, \tau) \frac{\partial^3 \psi}{\partial x^3} / \frac{1}{\psi(x, \tau)}$$

$$\frac{\partial \psi}{\partial \tau} \frac{\partial \psi}{\partial x} - \psi(x, \tau) \frac{\partial^2 \psi}{\partial x \partial \tau} = D \frac{\partial \psi}{\partial x} \cdot \frac{\partial^2 \psi}{\partial x^2} - D \psi(x, \tau) \frac{\partial^3 \psi}{\partial x^3}$$

→ nelineární III. stupně

$$\frac{\partial \psi}{\partial x} \left( \frac{\partial \psi}{\partial \tau} - D \frac{\partial^2 \psi}{\partial x^2} \right) - \psi(x, \tau) \frac{\partial}{\partial x} \left[ \frac{\partial \psi}{\partial \tau} - D \frac{\partial^2 \psi}{\partial x^2} \right] = 0$$

$$\frac{\partial}{\partial x} \left( \frac{\frac{\partial \psi}{\partial \tau} - D \frac{\partial^2 \psi}{\partial x^2}}{\psi(x, \tau)} \right) = 0$$

$$\frac{\partial \psi}{\partial \tau} - D \frac{\partial^2 \psi}{\partial x^2} = C \cdot \psi(x, \tau)$$

lin. PDE : lineární alternativa

Burgersova rovnice